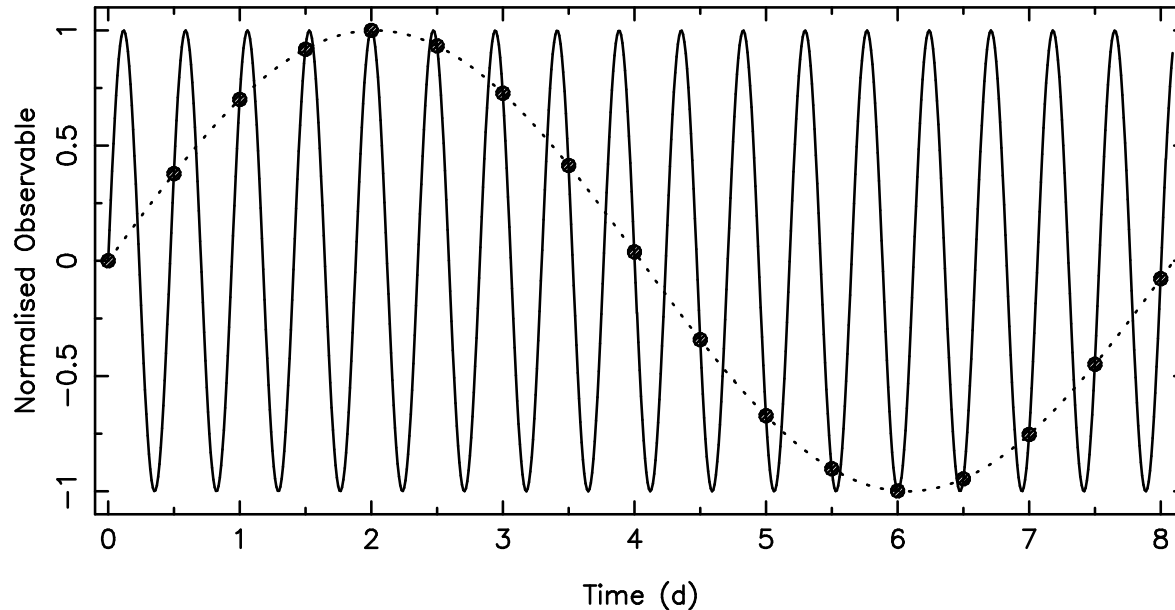


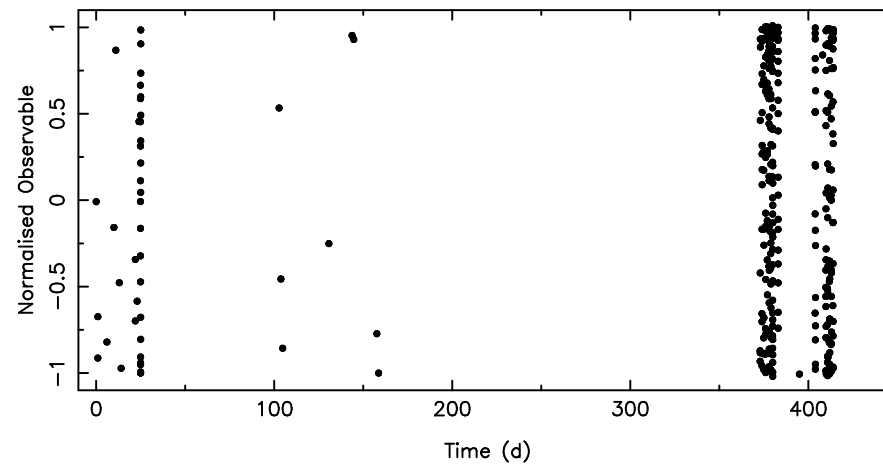
Introduction into Time Series Analysis of Unequally-spaced and Gapped Astronomical Data

1. Some preliminaries



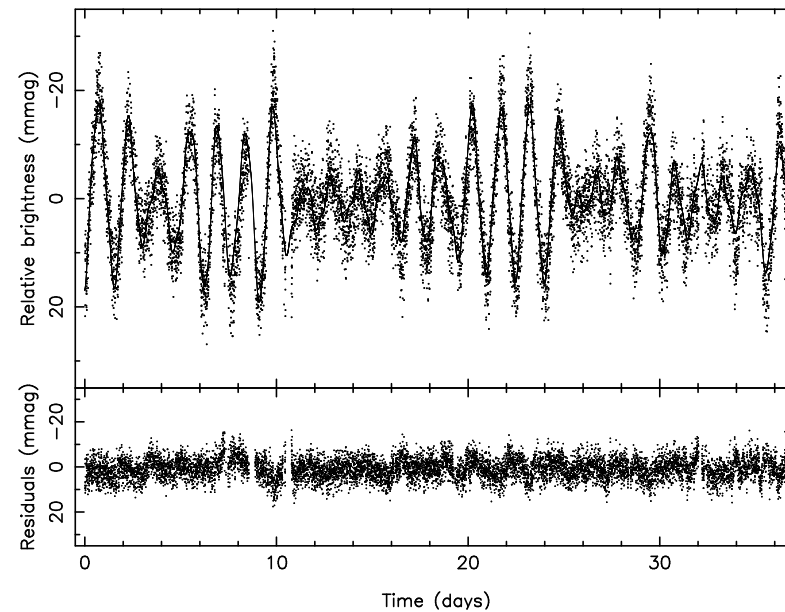
Simulated data (dots) representing a periodic signal with frequency $\nu = 0.123456789 \text{ d}^{-1}$. The dotted line is a harmonic fit for this frequency. The full line represents a fit with the frequency $2.123456789 \text{ d}^{-1}$.

1. Some preliminaries



Simulated gapped data representing a typical time series for a single-site campaign of a pulsating star.

1. Some preliminaries



Real data representing a typical time series for a space campaign of a pulsating star.

2. Methods based on Phase Dispersion Minimisation

Consider quantity x observed at t_i , $x_i(t_i)$ with $i = 1, \dots, N$.

The *phase* $\varphi(t_i)$ for *cyclic frequency* f or the *period* $P = 1/f$ with respect to the *reference epoch* t_0 is:

$$\varphi(t_i) = [f(t_i - t_0)] = \left[\frac{t_i - t_0}{P} \right],$$

with $0 \leq \varphi < 1$.

A plot of $x_i(t_i)$ as a function of $\varphi(t_i)$ is called a *phase diagram*.

For test period P : divide phase interval $[0, 1]$ in M equal *bins* with bin index $J_i = INT(M\varphi_i) + 1$, $INT(x) \equiv x - [x]$.

Suppose that the j -th bin contains N_j measurements.

Per bin:

$$\overline{x_j} = \sum_{i=1}^{N_j} \frac{x_{ij}}{N_j}, \quad V_j^2 = \sum_{i=1}^{N_j} (x_{ij} - \overline{x_j})^2 = \sum_{i=1}^{N_j} x_{ij}^2 - N_j \overline{x_j}^2,$$

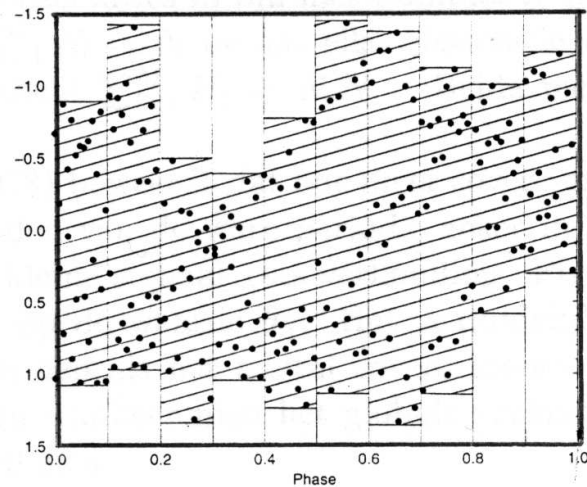
$$s_j^2 = \frac{V_j^2}{N_j - 1},$$

with x_{ij} the observation x_i with bin index $J_i = j$.

For whole dataset, \overline{x} , V^2 and s^2 , are defined as

$$\overline{x} = \sum_{i=1}^N \frac{x_i}{N}, \quad V^2 = \sum_{i=1}^N (x_i - \overline{x})^2 = \sum_{i=1}^N x_i^2 - N \overline{x}^2,$$

$$s^2 = \frac{V^2}{N - 1}.$$



For the M bins:

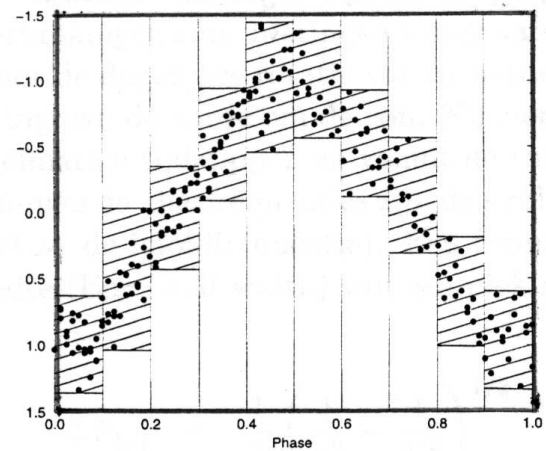
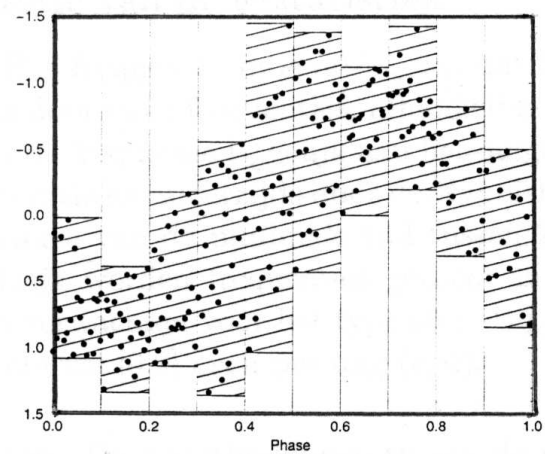
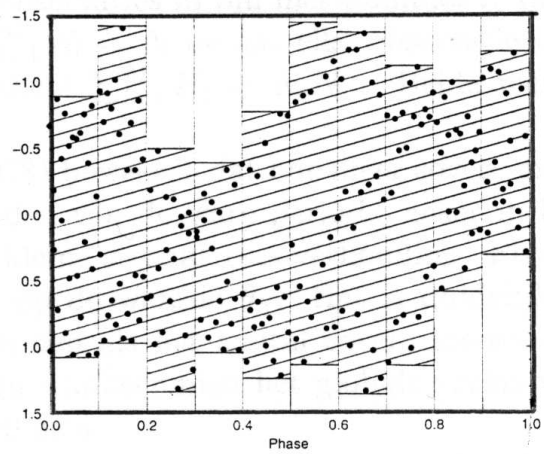
$$V_M^2 = \sum_{j=1}^M V_j^2, \quad V_G^2 = \sum_{j=1}^M N_j (\overline{x_j} - \overline{x})^2.$$

Thus $V^2 = V_M^2 + V_G^2$.

$\overline{x_j} \neq \overline{x}$ for good period: $V_G^2 \simeq V^2$

$\overline{x_j} \simeq \overline{x}$ for bad period: $V_G^2 \ll V^2$

search maximum of $V_G^2 \Leftrightarrow$ search minimum of V_M^2



Better: bin-cover structure (N_b, N_c) : divide phase diagram in N_b bins of length $1/N_b$ and apply this partition N_c times, with each partition shifted over $1/N_b N_c$

Define

$$\Theta \equiv \frac{\sum_{j=1}^M (N_j - 1) s_j^2 / \sum_{j=1}^M N_j - M}{\sum_{i=1}^N (x_i - \bar{x})^2 / N - 1} \in [0, 1]$$

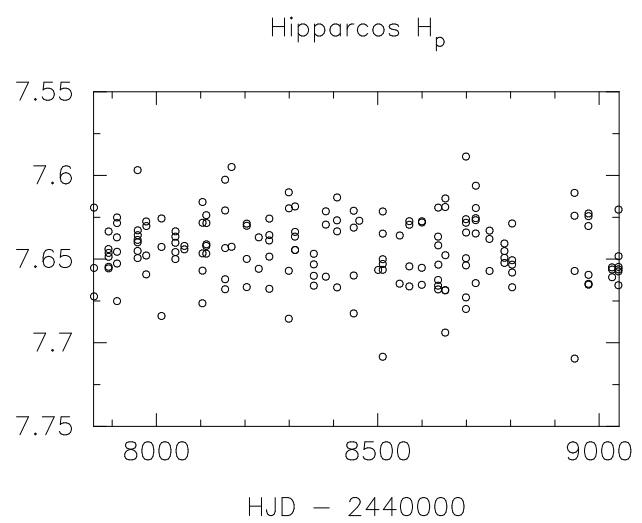
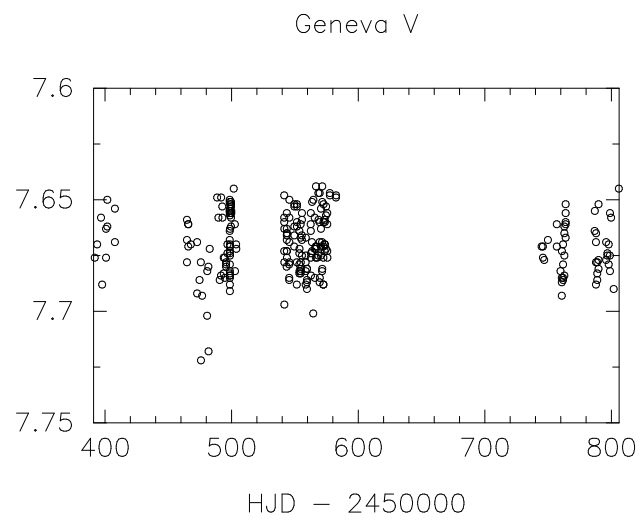
where s_j^2 is defined as:

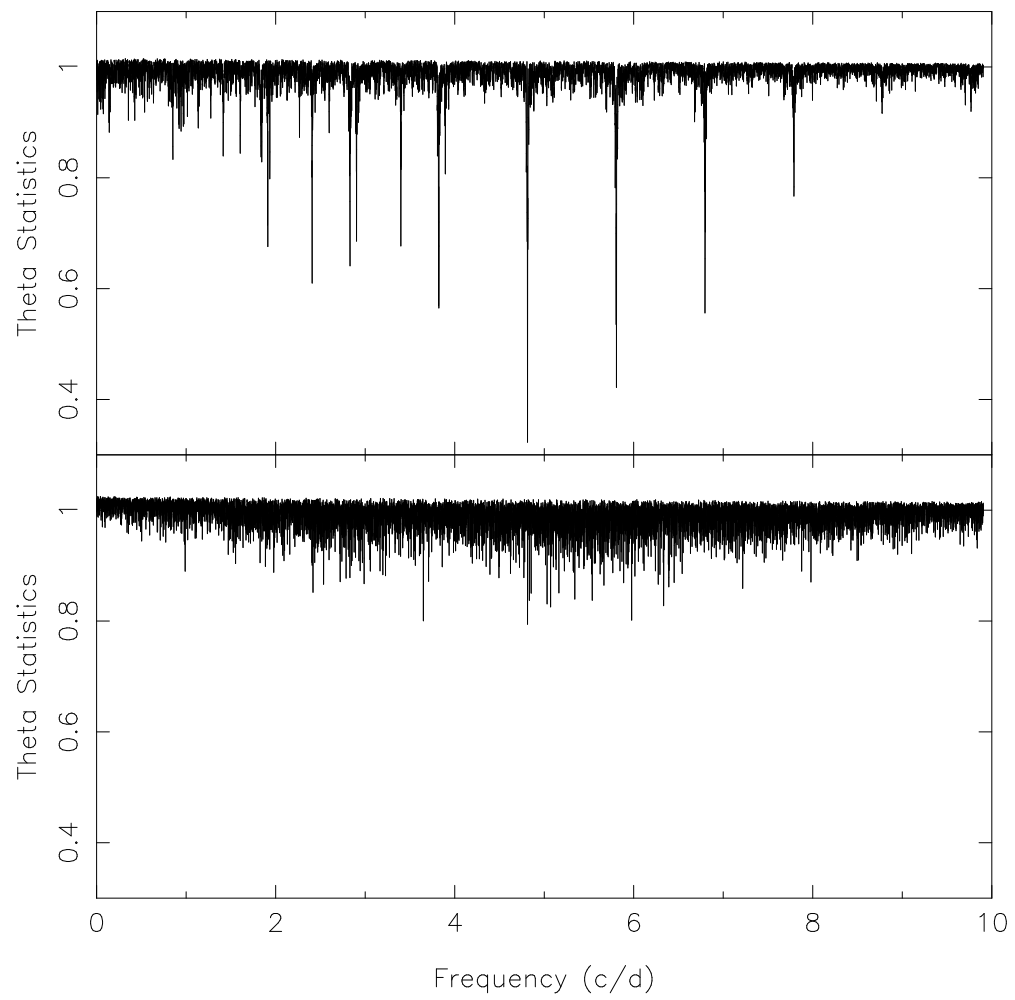
$$s_j^2 \equiv \frac{\sum_{i=1}^{N_j} (x_{ij} - \bar{x}_j)^2}{N_j - 1}.$$

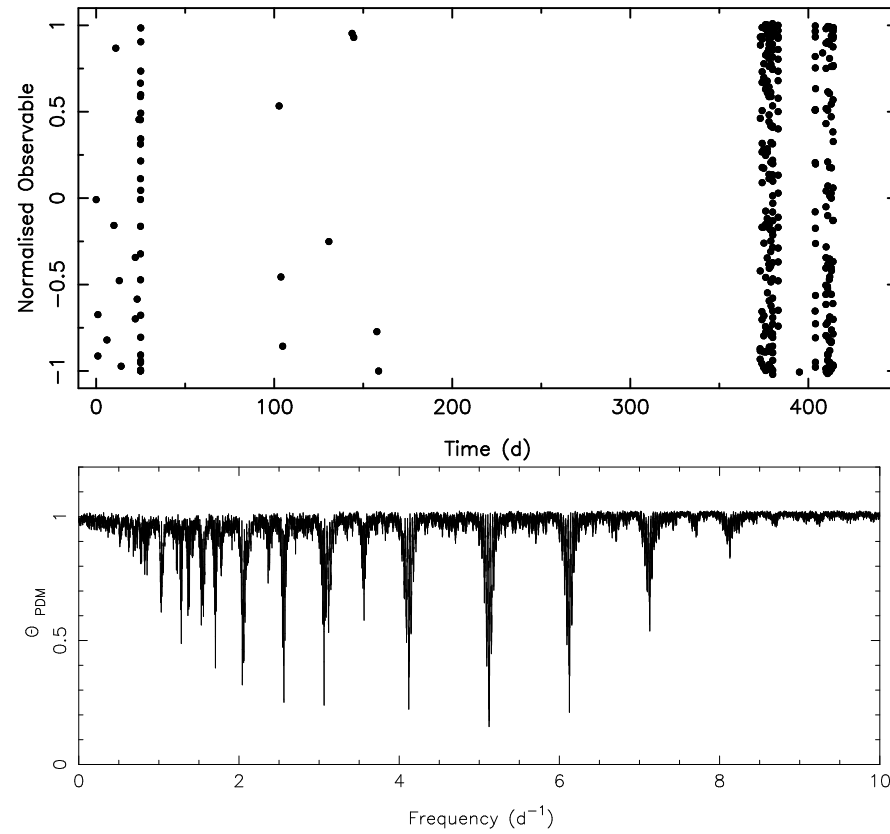
With the introduced notation:

$$\Theta = \frac{V_M^2 / \left(\sum_{j=1}^M N_j - M \right)}{V^2 / (N - 1)} = \frac{V_M^2 / N_c (N - N_b)}{V^2 / (N - 1)}.$$

A minimum in the Θ –statistic corresponds to a minimum of V_M

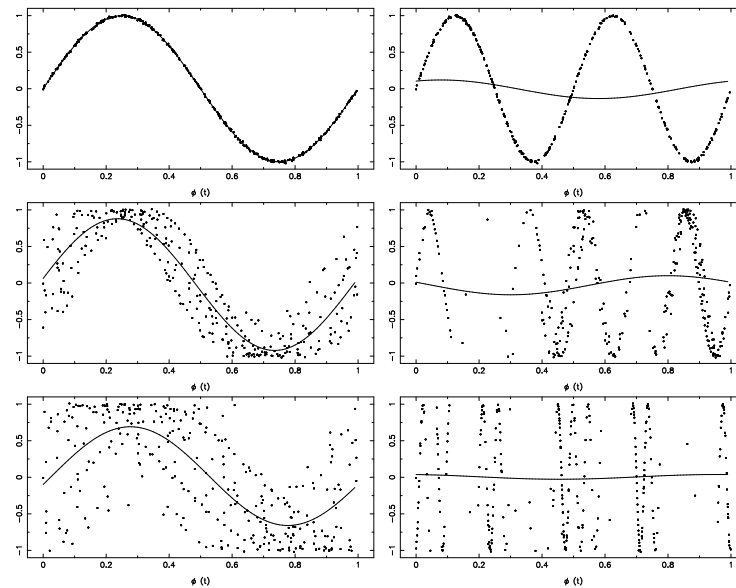






Statistic Θ_{PDM} of the data using 10 bins and 2 covers.

PDM and any other string-length method are sensitive to harmonics and subharmonics ... Be careful !



Phase diagrams for six minima in the Θ_{PDM} : 5.123 d^{-1} (upper left), 4.121 d^{-1} (middle left), 7.129 d^{-1} (lower left), 2.562 d^{-1} (upper right), 1.021 d^{-1} (middle right), 0.244 d^{-1} (lower right).

3. Methods based on Fourier Analysis

Fourier transform of $x(t)$:

$$F(f) \equiv \int_{-\infty}^{+\infty} x(t) \exp(2\pi i f t) dt$$

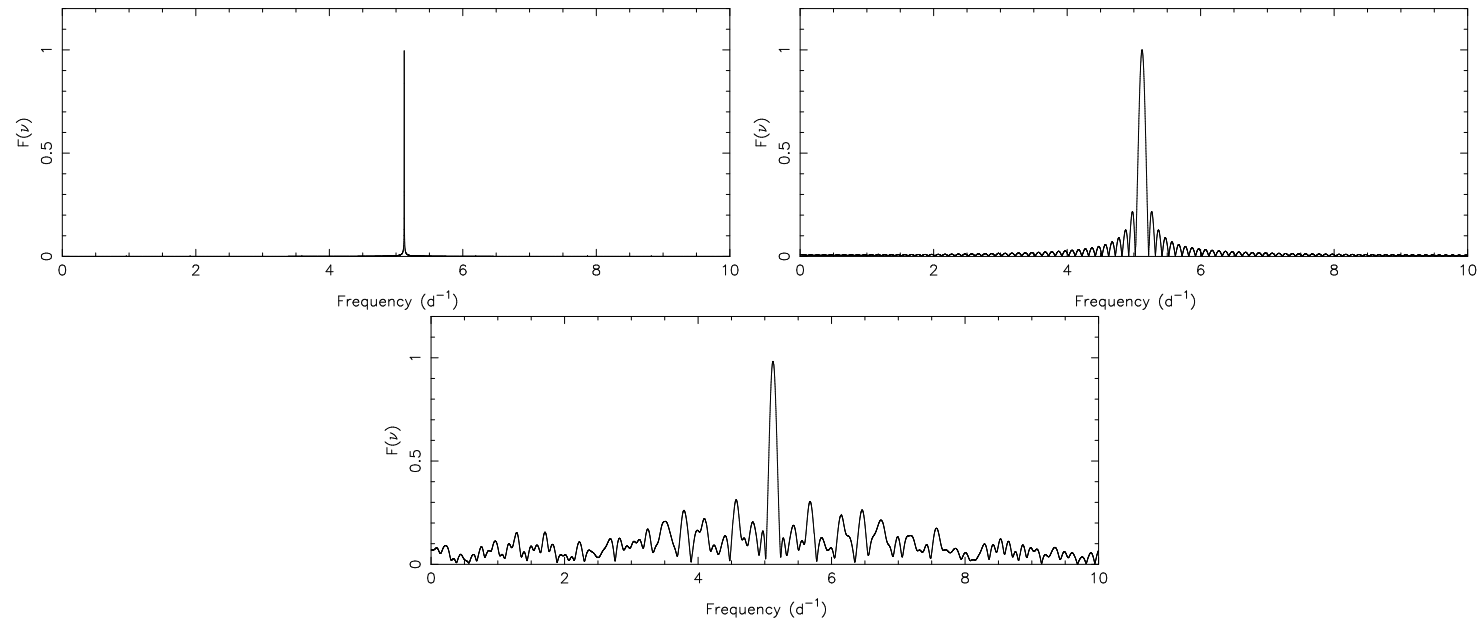
Fourier transform $F(f)$ of sum of harmonic functions with frequencies f_1, \dots, f_n and amplitudes A_1, \dots, A_n :

$$x(t) = \sum_{k=1}^n A_k \exp(2\pi i f_k t) : F(f) = \sum_{k=1}^n A_k \delta(f - f_k)$$

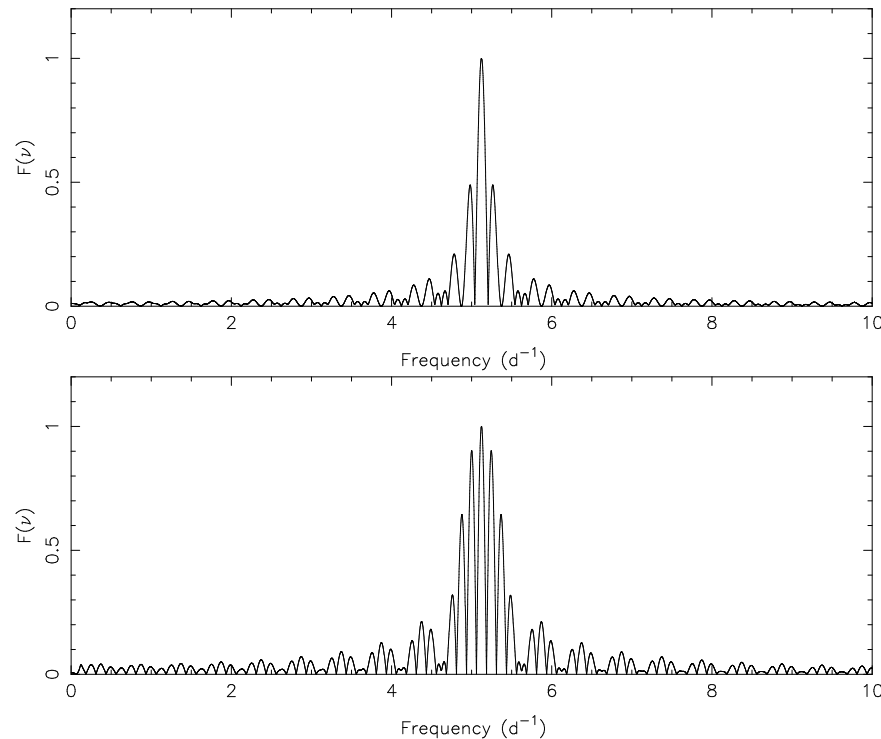
For $x(t) = \text{sine with frequency } f_1$, $F(f) \neq 0$ for $f = \pm f_1$

For $x(t) = \text{sum of } n \text{ harmonic functions with frequencies } f_1, \dots, f_n$,
 $F(f) = \text{sum of } \delta\text{-functions } \neq 0 \text{ for } \pm f_1, \dots, \pm f_n$

Real data set: $x(t)$ known for a discrete number of times $t_j, j = 1, \dots, N$



Left: 10^6 points over 1000 d, right: 10^4 points over 10 d, bottom: 4472 points over 10 d.



Fourier transforms of a noiseless time series of a sine function with frequency $5.123456789 \text{ d}^{-1}$ generated for a finite time span of 10 days and containing one large gap from day 4 until day 6 (top) and from day 2 until day 8 (bottom).

Discrete Fourier transform:

$$F_N(f) \equiv \sum_{j=1}^N x(t_j) \exp(2\pi i f t_j)$$

$F_N \neq F$! but connected through window function:

$$w_N(t) \equiv \frac{1}{N} \sum_{j=1}^N \delta(t - t_j)$$

Hence:

$$\frac{F_N}{N} = \int_{-\infty}^{+\infty} x(t) w_N(t) \exp(2\pi i f t) dt$$

Discrete Fourier transform of window function = spectral window $W_N(f)$:

$$W_N(f) = \frac{1}{N} \sum_{j=1}^N \exp(2\pi i f t_j)$$

Discrete Fourier transform = convolution of spectral window and Fourier transform: $F_N(f)/N = F(f) * W_N(f)$.

If $F(f) = \delta$ at f_1 , $F_N(f)/N$ same behaviour as $W_N(f)$ at f_1 :

$$F_N(f)/N = W_N(f) * \delta(f - f_1) = W_N(f - f_1)$$

\Downarrow

Comparison of $W_N(f)$ with $F_N(f)/N$ near $f_1 \rightarrow$ is f_1 is real or not?

If $F(f) =$ sum of n δ -functions:

$$\begin{aligned} \frac{F_N(f)}{N} &= W_N(f) * \sum_{k=1}^n \delta(f - f_k) = \sum_{k=1}^n W_N(f) * \delta(f - f_k) \\ &= \sum_{k=1}^n W_N(f - f_k) = \frac{1}{N} \sum_{k=1}^n \sum_{j=1}^N \exp(2\pi i (f - f_k)t_j). \end{aligned}$$

Hence, $F_N(f)/N =$ sum of n spectral windows centered around f_k

$W_N(f) \neq 0$ at all f (not only at $f_k, k = 1, \dots, n$) due to interference

\Downarrow

maxima in periodogram at spurious frequencies due to observing times. . .

ALIAS FREQUENCIES

Most common alias frequencies? Assume $t_j = t_0 + j\Delta t$:

$$\begin{aligned} W_N(f) &= \frac{1}{N} \sum_{j=1}^N \exp(2\pi i f t_0) \exp(2\pi i f j \Delta t) \\ &= \frac{1}{N} \exp(2\pi i f t_0) \sum_{j=1}^N \exp(2\pi i f j \Delta t) \\ &= \exp(2\pi i f t_0) \exp(\pi i f \Delta t (N + 1)) \frac{\sin(\pi f N \Delta t)}{N \sin(\pi f \Delta t)} \end{aligned}$$

in which we have made use of

$$\sum_{j=0}^{N-1} z^j = \frac{1 - z^N}{1 - z}$$

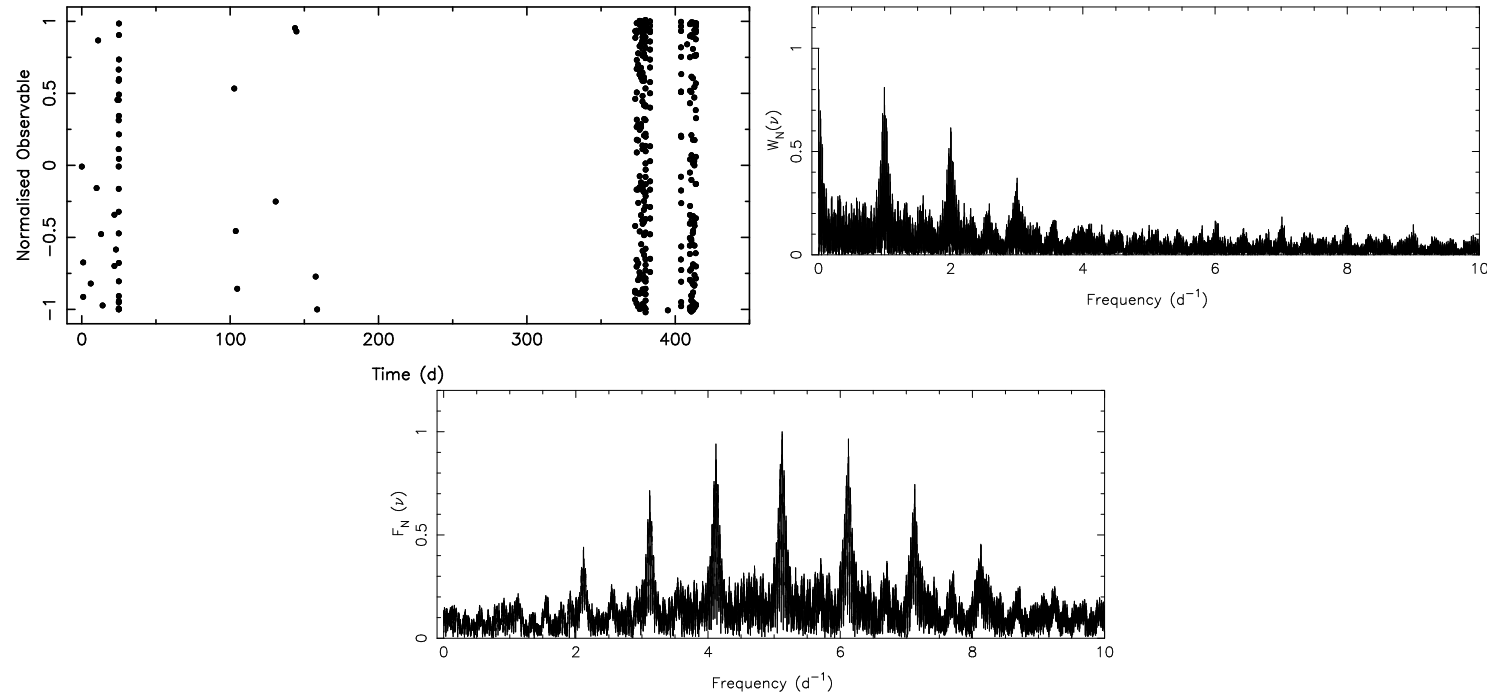
with $z = \exp(2\pi i f \Delta t)$.

For $t_0 = -(N + 1)\Delta t/2$: $W_N(f) = \frac{\sin(\pi N f \Delta t)}{N \sin(\pi f \Delta t)}$:
periodic function with period $1/\Delta t$:

$$\left| W_N \left(f + \frac{n}{\Delta t} \right) \right| = |W_N(f)|.$$

$F_N(f)$ has maxima at $f_n = n/\Delta t$

For non-equidistant data: $\Delta t \simeq 1$ day, 1 year, specific gaps,
length of nightly strings, etc.



Spectral window (upper right) and DFT (bottom) of a noise-free sinusoid with amplitude 1 at $5.123456789 \text{ d}^{-1}$ for the sampling shown upper left.

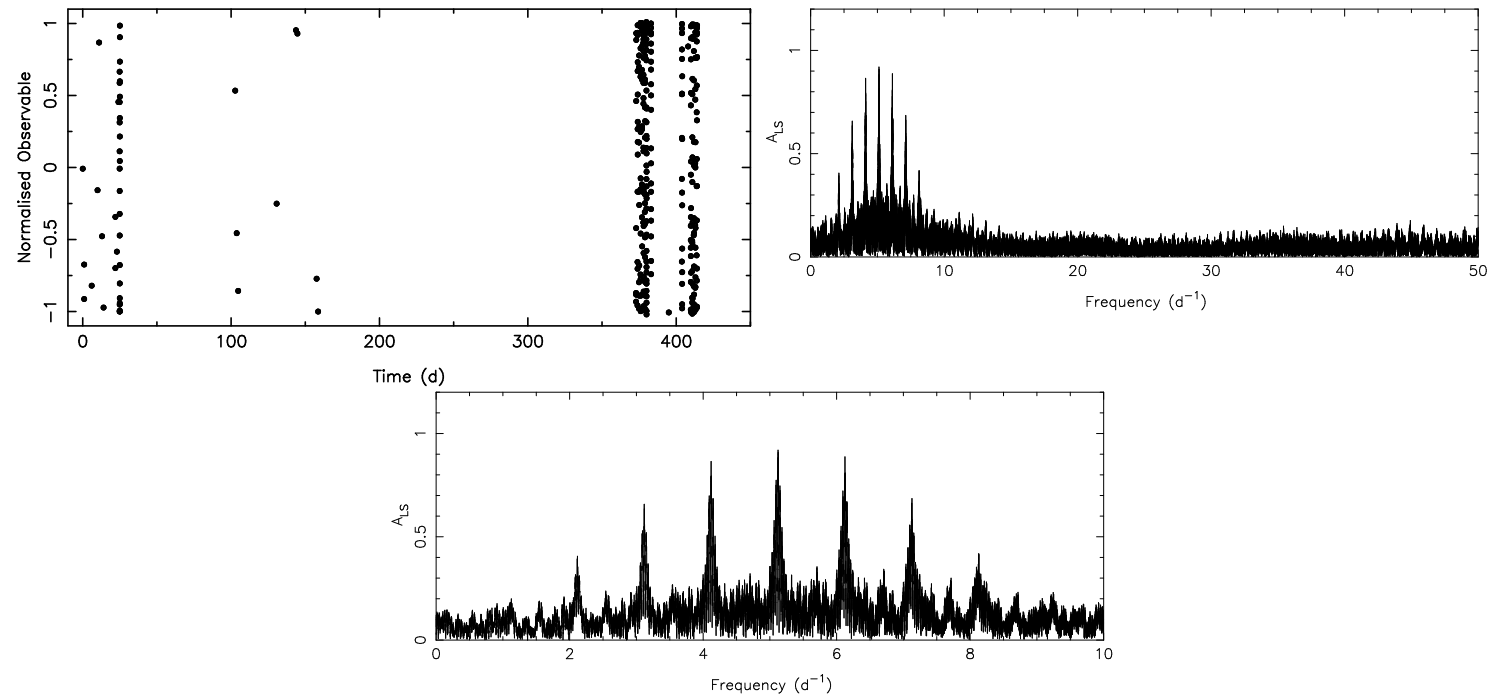
The Lomb-Scargle periodogram:

$$P_N(f) = \frac{1}{2} \left\{ \frac{\left\{ \sum_{i=1}^N x(t_i) \cos[2\pi f(t_i - \tau)] \right\}^2}{\sum_{i=1}^N \cos^2[2\pi f(t_i - \tau)]} + \frac{\left\{ \sum_{i=1}^N x(t_i) \sin[2\pi f(t_i - \tau)] \right\}^2}{\sum_{i=1}^N \sin^2[2\pi f(t_i - \tau)]} \right\}$$

with τ chosen such that

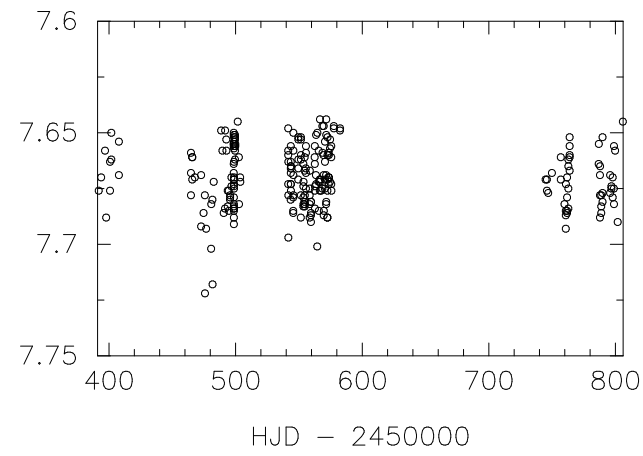
$$\sum_{i=1}^N \cos[2\pi f(t_i - \tau)] \sin[2\pi f(t_i - \tau)] = 0.$$

Note: $\lim_{N \rightarrow \infty} P_N(f) = A^2 N / 4$

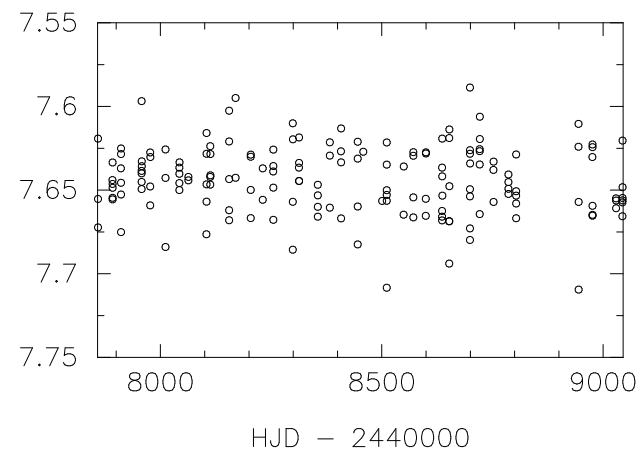


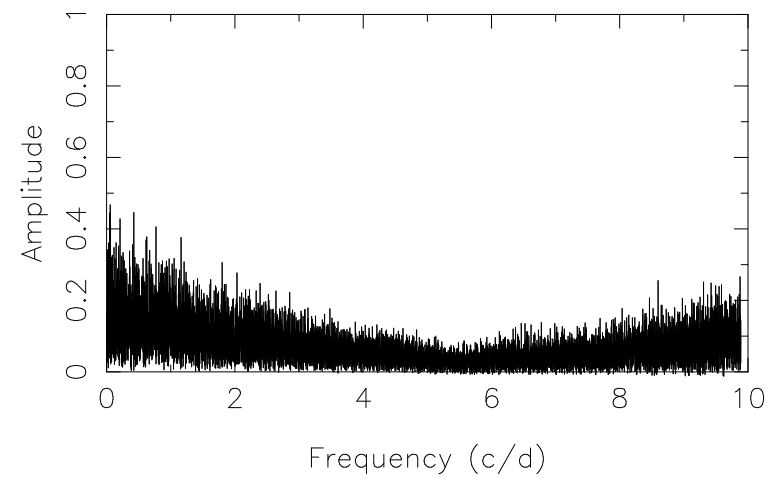
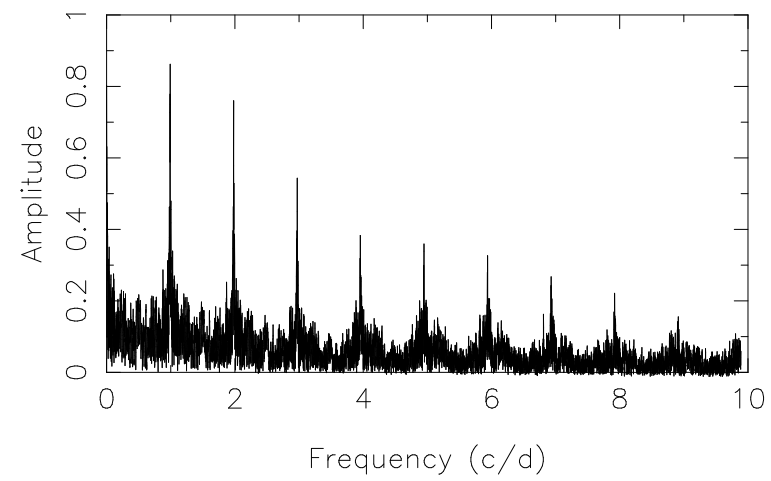
Lomb-Scargle periodograms for the data.

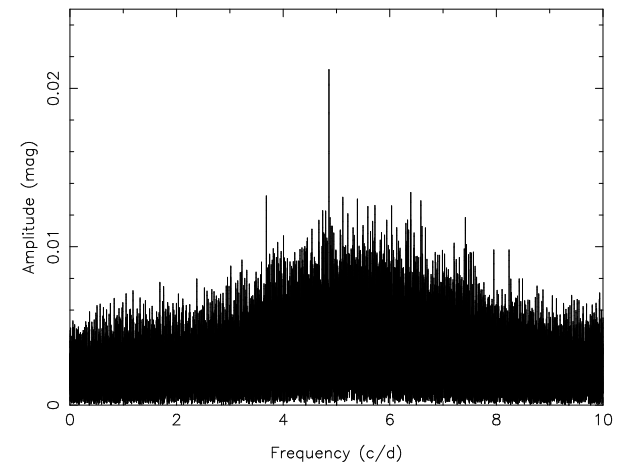
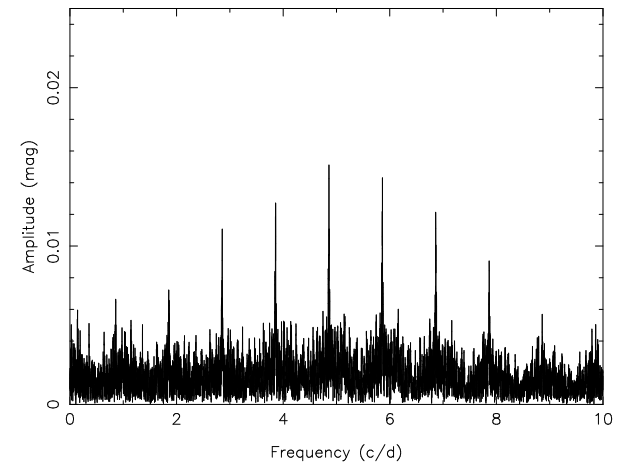
Geneva V



Hipparcos H_p







- Error estimate for derived frequency?

Assuming white noise and no alias problems:

$$\sigma_f = \frac{\sqrt{6}\sigma_R}{\pi\sqrt{N}AT}$$

with σ_R average error of data (\sim standard deviation of noise),
 T total time span, N number of data points

- Interval of testfrequencies?

use interval $[0, Nyq]$ with Nyq the Nyquist frequency.

For equidistant data: $1/2\Delta T$ with ΔT the time step.

For unevenly spaced data: take median of all time intervals between two consecutive measurements.

- Acceptance criterion for a frequency: $A > 4\sigma_{\text{Fourier}}$. Compute σ_{Fourier} in oversampled interval of testfrequencies around candidate frequency.

Criterion: 99.9% confidence interval that peak is not due to noise.

How good is the solution ?

LS fitting with f fixed:

$$x_i(t_i) = A \sin [2\pi(ft_i + \psi)] + C$$

Variance reduction $\in [0, 1]$:

$$1 - \frac{\sum_{i=1}^N \{x_i - [A \sin (2\pi(ft_i + \psi)) + C]\}^2}{\sum_{i=1}^N (x_i - \bar{x})^2}$$

Search for new frequencies in residuals $R_i(f) \equiv x_i - x_i^c(f)$ with:

$$x_i^c(f) \equiv A \sin [2\pi(ft_i + \psi)] + C$$

and so on, until amplitude $A < 4\sigma_{\text{Fourier}}$

At some point, you will reach the noise level ...

